Mechanochemical effects for some Al₂O₃ powders of dry grinding

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Three kinds of AI_2O_3 powders, i.e. two kinds of low-soda AI_2O_3 with average particle sizes of 3.9 and 0.6 µm and an electrofused AI_2O_3 with an average particle size of 21.8 µm, were ground for up to 300 h in a dry vibration ball mill. Variations in particle-size distribution, specific surface area, crystallite size, lattice strain, effective temperature factor and lattice constant were examined against milling time. The mechanism of grinding was found to differ between low-soda AI_2O_3 and electrofused AI_2O_3 . The mechanochemical effects on these AI_2O_3 powders occurred in the order decrease of crystallite size→increase of effective temperature factor \rightarrow increase of lattice strain. The length of the *a*-axis was clearly increased by a prolonged grinding. The difference in the ground state of three specimens was attributed to differences in the physical state of particles originating from the preparation methods, and also to particle size.

1. Introduction

Al₂O₃ powders have been widely used in the ceramics industry, and various kinds of powders: 1. sintered Al₂O₃; 2. low-soda Al₂O₃; 3. fused Al₂O₃; and 4. highpurity Al₂O₃, may be prepared by different methods [1]. Among these, the starting raw material for categories 1 and 3 is Al₂O₃ powder produced by the Bayer process. High-purity Al₂O₃ is prepared by various methods and has common properties such as a small (submicrometre) particle size and purity higher than 4N grade. On the other hand, the particle size of Al₂O₃ powders belonging to categories 1–3 is usually coarser than that of category 4. Therefore the grinding of these powders is considered to be one of the important essential processing steps.

It is well known that grinding treatment not only causes breakdown of particles but also has mechanochemical effects. The former factor corresponds to variations in particle size distribution and specific surface area. The latter appears as variation of crystallite size, lattice strain and effective temperature factor [2], which represents the degree of fluctuation of the atomic positions. Some work on the relation between grinding treatment and sinterability of Al₂O₃ powders has been reported [3-6]. Excess grinding causes a heavy agglomeration of the powders and also an increase in weight loss of the powders, and therefore was not suitable for sintering except for initial stage, which was enhanced by the activated surface diffusion. On the other hand, examination of the mechanochemical effects for Al₂O₃ powders has not yet been sufficiently elucidated.

Three kinds of Al_2O_3 powders, i.e. two kinds of lowsoda Al_2O_3 with different particle sizes and an electrofused AI_2O_3 were ground in a dry vibration ball mill. Variations in particle size, specific surface area, crystallite size, lattice strain, effective temperature factor and lattice constant were measured against milling time. From these data, the mechanisms of grinding and mechanochemical effects on the powders were investigated in relation to preparation methods.

2. Experimental procedure

The starting materials used for this study were two kinds of low-soda Al₂O₃ with an average particle size of 0.6 µm (LS23) and 3.9 µm (LS21), both of which are produced by Nippon Light Metal Co. Ltd, Shimidzu, Japan, and an electrofused Al₂O₃ with an average particle size of 21.8 µm (WA600) produced by Showa Denko Co. Ltd, Shiojiri, Japan. They reported chemical compositions are Al₂O₃ 99.8, Na₂O 0.02-0.06, SiO_2 0.02–0.04 and Fe_2O_3 0.03 wt % for the low-soda Al₂O₃; and Al₂O₃ 99.6, Na₂O 0.2, SiO₂ 0.02 and Fe_2O_3 0.03 wt % for the electrofused Al_2O_3 . Powder X-ray diffraction (XRD) revealed that the electrofused Al₂O₃ contained a small amount of β -Al₂O₃ besides a-Al₂O₃. Specimens LS23, LS21 and WA600 are designated S_{0.6}, S_{3.9} and S₂₂, respectively. To remove the influence of various pretreatments, powders were annealed at 800 °C for 2 h before the experiments. Dry vibration milling (MB-1, Chuoh Kakohki Co. Ltd) was performed up to 300 h with a vibration number of 1000 r.p.m. The Al₂O₃ powder of 700 g was charged in a 2-1 Al₂O₃ pot with 2 kg of Al₂O₃ balls (20 mm in diameter) as grinding media. No atmosphere control was used in the milling experiments. After a certain milling time, 15 g of the sample was collected from the

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pot and 1-4 g of grinding aid $(Ti(OC_4H_9)_4)$ was added each time to avoid agglomeration of the specimens [7] during milling. The accumulated amount of the TiO₂ component in the final powders was calculated to be 0.7 wt %.

Particle-size distribution was measured by a Microtrack instrument (Nikkisoh Co. Ltd) dispersing the powders ultrasonically in a HCl solution adjusted to pH 2. Specific surface area was measured by the BET method using a nitrogen gas with a Sorptomatic 1800 instrument (Carlo Elba Co. Ltd). Size and shape of particles were observed by SEM (MSM-9, Akashi Seisakusho Co. Ltd).

XRD measurement (Geigerflex, Rigaku Co. Ltd) was made with a step scanning and fixed-time method using graphite monochromated CuKa radiation. The experimental conditions were as follows: a slit system of $1-1^{\circ}-0.15-0.8$ mm, a step width of 0.004° , and a fixed time of 4 s. Peak position, peak intensity and integral width of twelve reflections from 20 to 140° in 20 were precisely measured and used for the calculation of lattice constant, effective temperature factor, crystallite size and lattice strain, respectively. Since the reflections of higher angles were heavily diffused and weakened by milling in specimen $S_{0.6}$, only six reflections could be used for the various calculations. The standards for correction of instrumental broadening of reflections and that of reflection angles were the specimen S_{22} annealed at 1600 °C for 2 h and the Si powder (640a, NBS), respectively. Crystallite size (D_c) and lattice strain (η) were calculated by Hall's method [8] using the following formula:

$$\beta \cos \theta / \lambda = 1/D_{c} + 2\eta \sin \theta / \lambda \qquad (1)$$

where β is a real integral width of specimen after corrected for instrumental broadening, and λ is the X-ray wavelength. Effective temperature factor (B_{eff}) was calculated by the following formula [2]:

$$\ln (I_{\rm obs}/I_{\rm std}) = \ln k - 2B_{\rm eff}(\sin^2\theta/\lambda^2) \qquad (2)$$

where I_{obs} and I_{std} were the observed intensity of the specimen and that of the standard, and k is a constant. Lattice constants of α -Al₂O₃ were calculated by the least-squares method using the RSLC-3 program [9].

3. Results

3.1. Particle size and specific surface area

Fig. 1 shows SEM photographs of as-received specimens. Particles of specimen S_{22} were an irregular and angular shape. Most of the particles of specimen $S_{3,9}$ were thick hexagonal plates. Particles of $S_{0.6}$ were small and rounded.

Fig. 2 shows variations of specific surface area for three specimens against milling time. The increase of specific surface area of S_{22} was much higher than those of specimens $S_{0.6}$ and $S_{3.9}$. This difference must correspond to the difference in grinding mechanism. It was also apparent in the difference in curve shapes that the curves of $S_{0.6}$ and S_{22} were convex, but that of $S_{3.9}$ was concave.

Fig. 3 shows the variation in average particle size (d_{50}) , a spherically corresponding particle size estim-



Figure 1 SEM photographs of as-received specimens. (a) S_{22} ; (b) $S_{3,9}$; (c) $S_{0,6}$.

ated from the specific surface area (d_s) , and a crystallite size (D_c) for three specimens against milling time. The d_{50} of the three specimens decreased with milling time: $23.8 \rightarrow 3.4 \ \mu\text{m}$ in S_{22} , $5.2 \rightarrow 3.5 \ \mu\text{m}$ in $S_{3.9}$, and $1.35 \rightarrow 0.97 \ \mu m$ in S_{0.6}. (Here, since specimens S₃₀₉ and $S_{0.6}$ after annealing were heavily agglomerated by the annealing treatment, the d_{50} data after 20 h milling were used for the initial data.) Variation of the d_{50} in $S_{0.6}$ and $S_{3.9}$ was apparently smaller than that in S_{22} . The d_{50} of specimen S_{0.6} decreased a little up to 180 h milling time, but was almost constant with further milling time. Variations in d_s showed a similar trend to those of d_{50} . The d_s curves of S_{0.6} and S₂₂ showed a concave shape and it was considered to be the tendency for saturation. That of specimen $S_{3,9}$, however, showed a constant decrease up to 300 h milling.

3.2. Mechanochemical parameters

Variation of crystallite size for three specimens against milling time is shown in Fig. 3. Crystallite sizes of asannealed specimens $S_{3,9}$ and S_{22} were very large and



Figure 2 Variation of specific surface area against milling time for each specimen. \bigcirc , $S_{0.6}$; \triangle , $S_{3.9}$; \blacklozenge , S_{22} .

were taken as infinite size by this method. They decreased very rapidly as soon as the milling started and decreased to 100 nm in $S_{3.9}$ and 88 nm in S_{22} after 300 h milling. On the other hand, the crystallite size of $S_{0.6}$ was evidently smaller than those of the other two specimens even after annealing; it decreased further from 140 to 77 nm with 100 h milling, and kept a constant value even with further milling.

Fig. 4 shows variation in lattice strain for three specimens against milling time. The lattice strains of three as-annealed specimens showed a little difference, and that of S_{22} was the largest. Thermal stress due to the very high temperature process for the preparation of this specimen was considered to be the reason. The lattice strain was almost constant during the initial milling period up to 40 h because mechanochemical energy was almost expended by the rapid decrease in crystallite size in this period. After this period, the lattice strain began to increase with milling time. The increase in lattice strain in $S_{3,9}$ was the smallest, but maintained a constant increasing rate even after 300 h milling. The other two specimens showed a convex curve and the increase in lattice strain in S_{22} seemed to be almost saturated by this milling time.

Fig. 5 shows variation in the effective temperature factor for three specimens against milling time. The effective temperature factors increased with the milling time, but their rates were not a constant against milling time. An increasing rate between around 50 and 100 h was especially high. The effective temperature factors of S_{22} and $S_{0.6}$ became almost similar values after 300 h milling, and seemed to be almost saturated. That of $S_{3.9}$ was, however, only about half of those of the other specimens, but seemed to continue to increase linearly even after 300 h milling.

Fig. 6 shows the variation of the lattice constant of α -Al₂O₃ against milling time. The longer the milling time, the length of the *a*-axis showed a more significant increase up to 220 h milling, while no change or a small decrease was observed for the length of the *c*-axis. Here the *a*-axis is parallel to the layer of the h c p structure, and the *c*-axis is perpendicular to it.



Figure 3 Variation of average particle size (d_{50}) , spherically corresponding particle size estimated from specific surface area (d_s) , and crystallite size (D_c) against milling time for each specimen. Circles, S_{0.6}; triangles, S_{3.9}; diamonds, S₂₂.



Figure 4 Variation of lattice strain (η) against milling time for each specimen. \bigcirc , S_{0.6}; \triangle , S₃₉; \blacklozenge , S₂₂.



Figure 5 Variation of effective temperature factor (B_{eff}) against milling time for each specimen. \bigcirc , S_{0.6}; \triangle , S_{3.9}; \blacklozenge , S₂₂.

Since the values of a and c axes for as-annealed specimens were in a good agreement with those reported for α -Al₂O₃ [10], variations in lattice length, especially in the a-axis, were considered to be uniform lattice deformation caused by the grinding. It was not certain what kind of structural change corresponded to this lattice deformation. Considering that the deformation in the a-axis was much larger than that in the c-axis, introduction of a point defect or disordering of atomic position in the layer structure by sliding and/or atomic disordering may be plausible. Uniform lattice deformation by grinding has been reported in Mg(OH)₂ [11]; however expansion of c-axis, which is perpendicular to the layer structure, was observed in this case.



Figure 6 Variation of the lattice constant against milling time. Error bars represent $\pm 3\sigma$ ($\sigma =$ standard deviation) range. The dashed lines show reported data for α -Al₂O₃ [10]. \bigcirc , S_{0.6}; \triangle , S_{3.9}; \blacklozenge , S₂₂.

4. Discussion

The relation between applied grinding energy, E, and particle size, R, has been examined by many workers. The general formula to represent this relation can be written as follows [12]:

$$dE = kR^{-n}dR \qquad (3)$$

where k is a constant. Rittinger [13], considering the grinding mechanism, suggested that applied energy is almost expended by the increase in surface area of the particles. Assuming that no particle shape change occurs due to the grinding, Equation 3 can be integrated to the following formula:

$$E = k(S - S_0) = k'(1/R - 1/R_0)$$
(4)

where S is surface area; S_0 is initial surface area; R_0 is initial particle size; and k and k' are constants. This formula corresponds to an integration of Equation 3 by n = 2.

On the other hand, Kick [13] proposed a new theory relating applied grinding energy to volume change of particles. The following equation can be derived from this consideration, and corresponds to an integration of Equation 3 by n = 1.

$$E = k \ln (R_0/R) \tag{5}$$

Bond [14], considering the grinding mechanism, suggested that applied grinding energy is proportional to the variation of particle size, and proposed the intermediate multiplier n = 1.5 for the integration of Equation 3. Thus Equation 6 can be derived.

$$E = k(1/R^{1/2} - 1/R_0^{1/2})$$
 (6)

Assuming that the applied grinding energy is proportional to milling time, we attempted to determine the most probable grinding mechanism for three Al_2O_3 powders among the above three theories, using Equations 4–6. From the relation between specific surface area and milling time, the data of specimen S_{22} were found to be best fitted with Equation 4; however the data for the other specimens clearly did not fit with this formula. On the other hand, the data for $S_{3,9}$ and $S_{0.6}$ showed a good fit with Equation 5. It was therefore concluded that the grinding mechanism for electrofused Al_2O_3 and low-soda Al_2O_3 was different. This difference was thought to be produced by the difference in physical state of the particles. The particles of electrofused Al_2O_3 were considered to be composed with polygrains, and those of low-soda Al_2O_3 were single-grain, except for agglomerated particles. Therefore breakdown of the particles easily occurred from the grain boundary within a particle in electrofused Al_2O_3 , but not in low-soda Al_2O_3 .

The decrease of particle size in $S_{3.9}$ and $S_{0.6}$ was very small even after prolonged grinding. The high Young's modulus and hardness of Al_2O_3 were thought to be likely reasons. The large difference between the sizes of grinding balls and the particles of the specimens may be another reason. Grinding of $S_{0.6}$ under these conditions, therefore, seemed to reach a grinding limit [15], judging from the data for specific surface area change in this specimen.

The major mechanochemical energy on the particles appeared as variation in crystallite size, lattice deformation, effective temperature factor, and lattice strain. Here the crystallite size corresponds with domain size of the single-crystal region. The lattice deformation means a small expansion and/or shrinkage of the unit cell. Impurities, point defects and distortion of structure were considered for these reasons. The effective temperature factor represents the degree of vibration of atomic position and corresponds to randomization of a structure. The ultimate state is, therefore, an amorphous structure. The lattice strain corresponds to the distribution of compressively and tensilely distorted lattices, and the average lattice size was unchanged in this case. It often originated from the evolution of dislocations.

Fig. 7 shows the relation between crystallite size and lattice strain for each specimen. In the initial grinding stage, mechanochemical energy was solely expended by the decrease in crystallite size of the specimens. After decreasing crystallite size, the lattice strain began to increase. Since the crystallite size of $S_{0.6}$ was very small from the starting state, the lattice strain increased from the initial stage in this specimen. This crystallite size of 75 nm was considered to be the minimum under these grinding conditions. The lattice strain of S_{22} was larger than that of $S_{3,9}$ when comparing these for the same crystallite size. The higher temperature process applied to the former specimen probably causes this difference. Fig. 8 shows the relation between crystallite size and effective temperature factor for each specimen, and shows a similar trend to Fig. 7. The difference in effective temperature factor of S_{22} and that of $S_{3.9}$ was clearer than that in lattice strain shown in Fig. 7. This effect was again attributed to the higher temperature process for this specimen. The effective temperature factor of S_{22} became a closely similar value with that of $S_{0.6}$ after prolonged grinding, and seemed to be almost saturated by 300 h milling.



Figure 7 Relation between crystallite size and lattice strain for each specimen. \bigcirc , S_{0.6}; \triangle , S_{3.9}; \blacklozenge , S₂₂.



Figure 8 Relation between crystallite size and effective temperature factor for each specimen. \bigcirc , S_{0.6}; \triangle , S_{3.9}; \blacklozenge , S₂₂.

Fig. 9 shows the relation between lattice strain and effective temperature factor for each specimen. These data can be represented by three divided straight lines. An interesting point was that the data for $S_{3,9}$ and $S_{0,6}$ were plotted on the same straight lines. Initially, the increase in effective temperature factor was much more dominant than that for the lattice strain. Both effective temperature factor and lattice strain increased during the next stage. Finally, the effective temperature factor became saturated but the lattice



Figure 9 Relation between lattice strain and effective temperature factor for each specimen. \bigcirc , $S_{0.6}$; \triangle , $S_{3.9}$; \blacklozenge , S_{22} .

strain still increased. The data for S22 showed a similar trend to low-soda Al_2O_3 , but the data were plotted on different lines. Some difference was observed between the two types of Al₂O₃ specimens, i.e. low-soda and electrofused Al₂O₃. One difference was that the increase in effective temperature factor of electrofused Al_2O_3 during the initial stage was more noticeable than that in low-soda Al_2O_3 . This means that vibration of atomic position occurred more easily in the electrofused Al_2O_3 than in the low-soda Al_2O_3 . A second point is that a residue lattice strain of around 0.04% was observed in the electrofused Al_2O_3 when the data were extrapolated to $B_{eff} = 0$, while no residue lattice strain was observed in the low-soda Al₂O₃. It was considered from this result that annealing for the electrofused Al₂O₃ was insufficient under these conditions (e.g. 800 °C for 2 h) although for low-soda Al_2O_3 the same annealing conditions were sufficient.

Inagaki *et al.* [2] examined the relations between effective temperature factor and lattice strain for the various materials, and found that the effective temperature factor of most of the materials became saturated by prolonged grinding, but the lattice strain continued to increase. The effective temperature factor was considered to represent the degree of rigidity of the structure. Inagaki *et al.* [2] reported the minimum value for perovskite structure and the maximum value for graphite. Comparing the results for Al_2O_3 , this was ranked in the middle class among the various materials, and had a similar value to TiO_2 .

The mechanochemical effect for these specimens can be summarized as follows:

1st stage: decrease in crystallite size;

2nd stage: increase in effective temperature factor; 3rd stage: increase in both effective temperature factor and lattice strain; 4th stage: saturation of uniform lattice deformation and effective temperature factor, but increase in lattice strain.

This order of mechanochemical factors was considered to reflect the order of susceptibility of deformation. Variation in these factors therefore saturated in the same order. Flexible structures such as graphite [2], molecular bonded layer structures, and SiO₂ [16], constructed by the framework of corner-shared SiO_4 tetrahedra, easily convert to amorphous structures by prolonged grinding and their effective temperature factors can become very large. On the other hand, it was soon saturated in a rigid structure and the mechanochemical energy was stored by the lattice strain instead of the effective temperature factor. It was therefore considered that Al₂O₃ follows the latter case. After 300 h milling, the ground state of three Al_2O_3 specimens reached the following stages: 2nd stage in $S_{3,9}$; 3rd stage in S_{22} ; and 4th stage in $S_{0,6}$. These differences were derived from the difference in the physical state of the particles, i.e. polygrain or single grain, and also the particle size of the specimens.

5. Conclusions

Three kinds of Al_2O_3 powders, i.e. two kinds of lowsoda Al_2O_3 with average particle sizes of 3.9 and 0.6 µm and an electrofused Al_2O_3 with 21.8 µm, were ground by a dry vibration ball mill. The effect of grinding mechanism and the mechanochemical effects on the powders were investigated. To summarize, the following conclusions can be drawn.

1. The grinding mechanisms for Al_2O_3 powders differed with their preparation methods. Variation of particle size, R, with applied grinding energy was proportional to $\ln R$ in the low-soda Al_2O_3 and to 1/R in the electrofused Al_2O_3 .

2. The mechanochemical effects on the powders occurred in the order: decrease in crystallite size \rightarrow increase in effective temperature factor \rightarrow increase in lattice strain.

3. The increase in the effective temperature factor of electrofused Al_2O_3 occurred much more easily than for low-soda Al_2O_3 .

4. The effective temperature factor reached a constant value of around 2×10^{-2} nm⁻², and was in the middle class among various inorganic materials.

5. Uniform lattice distortion such as lengthening of the *a*-axis in α -Al₂O₃ was observed with grinding.

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